# Singular half-sided modular inclusions and deformation quantization 

Gandalf Lechner<br>joint work with Charley Scotford

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## Overview of the talk

- Main topic: Certain (half-sided modular, HSM) inclusions $\mathcal{N} \subset \mathcal{M}$ (subfactors of type $\mathrm{III}_{1}$ ) of von Neumann algebras (equivalently formulated as Borchers triples)
- Closely related to chiral conformal quantum field theories on the real line $\mathbb{R}$ (or the circle $S^{1} \cong \mathbb{R} \cup\{\infty\}$ )
- This relation will guide us to the question whether a $\mathrm{HSM} \mathcal{N} \subset \mathcal{M}$ is singular (trivial relative commutant, $\mathcal{N}^{\prime} \cap \mathcal{M}=\mathbb{C}$ ) or not.
- Will construct examples of singular half-sided inclusions by a deformation procedure (warped convolution, Rieffel deformation).


## Geometric preliminaries

- Symmetries of real line $\mathbb{R}$ :
- translations $x \mapsto x+a$, dilations $x \mapsto e^{t} \cdot x$, reflection $x \mapsto-x$
- generate affine group Aff of $\mathbb{R}$
- In conformal theories, have also conformal (Möbius) transformations

$$
x \mapsto \frac{a x+b}{c x+d}, \quad x \in \mathbb{R} \cup\{\infty\} \cong S^{1}
$$

generate Möbius group Möb $=\operatorname{PSL}(2, \mathbb{R})$

## Conformal nets

Idea: Consider unitary rep $U$ of Möb on Hilbert space $\mathcal{H}$ and model quantum fields localized in interval $I$ by an algebra of operators $\mathcal{A}(I) \subset \mathcal{B}(\mathcal{H})$.

## Definition (or wishlist:)

- For every $I \in \mathcal{I}$, there is a (von Neumann) algebra $\mathcal{A}(I) \subset \mathcal{B}(\mathcal{H})$
- $I_{1} \subset I_{2} \Longrightarrow \mathcal{A}\left(I_{1}\right) \subset \mathcal{A}\left(I_{2}\right)$
- Locality: If $I_{1}$ and $I_{2}$ are disjoint, then $\mathcal{A}\left(I_{1}\right)$ and $\mathcal{A}\left(I_{2}\right)$ commute.
- Covariance: $U(g) \mathcal{A}(I) U(g)^{-1}=\mathcal{A}(g I)$
- Positivity: The translations $U(x)=e^{i P x}$ have positive generator $P>0$
- Vacuum I: There is a unique (up to a scalar) vector $\Omega \in \mathcal{H}$ that is invariant under $U$.
- Vacuum II: $\Omega$ is cyclic for all $\mathcal{A}(I)$

Such a structure is called a conformal net.

- Although well motivated, it is a complicated definition. For a classification, one would like to connect it to simpler data.


## Borchers triples and half-sided inclusions

Let us start from a simpler situation (on the real line):

- only one algebra, the one corresponding to the half line $\mathbb{R}_{+}$
- only translation symmetry


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A (one-dimensional) Borchers triple $(\mathcal{M}, T, \Omega)$ consists of a von Neumann algebra $\mathcal{M} \subset \mathcal{B}(\mathcal{H})$ and a unitary representation $T$ of $\mathbb{R}$ on $\mathcal{H}$ s.t.

- $T$ has positive generator. The subspace of $T$-invariant vectors is $\mathbb{C} \Omega$.
- $T(x) \mathcal{M} T(-x) \subset \mathcal{M}$ for $x \geq 0$.
- $\Omega$ is cyclic and separating for $\mathcal{M}$.


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- $\Omega$ is cyclic and separating for $\mathcal{M}$.
- Gives rise to half-sided modular inclusion (Borchers, Wiesbrock)

$$
\begin{array}{r}
\mathcal{N} \subset \mathcal{M}, \quad \mathcal{N}:=T(1) \mathcal{M} T(-1), \\
\Delta_{\mathcal{M}}^{i t} \mathcal{N} \Delta_{\mathcal{M}}^{-i t} \subset \mathcal{N} \quad t \leq 0
\end{array}
$$

- To construct a map (Borchers triples) $\rightarrow$ (conformal nets), one needs modular theory.


## The modular symmetry machine

With modular theory, we can extend a Borchers triple to a conformal net [Borchers, Wiesbrock, Longo/Guido/Wiesbrock]

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- Define dilation symmetry by modular unitaries $\Delta_{\mathcal{M}}^{i t}$ of $(\mathcal{M}, \Omega)$.
- Define interval algebras by

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\mathcal{A}(a, b):=T(a) \mathcal{M} T(-a) \cap T(b) \mathcal{M}^{\prime} T(-b)
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- If $\Omega$ is cyclic for the interval algebras $\mathcal{A}(a, b)$ ("standard" situation), use modular unitaries $\Delta_{\mathcal{A}(a, b)}^{i t}$ to generate representation of Möbius group.


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## Theorem ([Longo,Guido, Wiesbrock 98])

- In the standard situation, this construction yields a conformal net on $S^{1}$.
- There exists a bijection between (strongly additive) conformal nets and standard Borchers triples.


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Unfortunately, the "standard" situation is not really standard ..

## The local subspace

Define the local subspace:

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\mathcal{H}_{\text {loc }}:=\overline{\mathcal{A}(I) \Omega} \subset \mathcal{H}, \quad I \subset \mathbb{R} \text { bounded interval }
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Theorem ([Bostelmann, GL, Morsella 11])
This space is independent of $I$ and invariant under the net $\mathcal{A}$.
Three cases:
(1) $\mathcal{H}_{\text {loc }}=\mathcal{H}$. (standard situation)

- Here we can construct a conformal net directly on $\mathcal{H}$.


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(2) $\mathbb{C} \Omega \mp \mathcal{H}_{\text {loc }}^{\ddagger} \mathcal{H}$.
- Here the construction works as in (1) after restriction to $\mathcal{H}_{\text {loc }}$.
(3) $\mathcal{H}_{\text {loc }}=\mathbb{C} \Omega$.
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Does case (3) occur?

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Does case (3) occur? Answer from 2019: Yes. [Longo, Tanimoto, Ueda 19] have free probability construction to get an example of (3).

## The algebra at infinity

Let $(\mathcal{M}, T, \Omega)$ be a BT, write $\alpha_{x}=\operatorname{Ad} T(x), \sigma_{t}=\operatorname{Ad} \Delta^{i t}, \mathcal{N}=\alpha_{1}(\mathcal{M})$.
The algebra at infinity:

$$
\mathscr{X}:=\bigcap_{t \in \mathbb{R}} \sigma_{t}(\mathcal{N} \vee J \mathcal{N} J)=\bigcap_{I \in \mathcal{I}} \mathcal{A}(I)^{\prime} .
$$

Remarks/Lemmas:

- The larger $\mathscr{X}$, the smaller $\mathcal{H}_{\text {loc }}$.
- $\mathscr{X}=\mathbb{C} 1 \Longleftrightarrow \mathcal{H}_{\text {loc }}=\mathcal{H}$ (standard case).
- $\mathscr{X}=\mathcal{B}(\mathcal{H}) \Longleftrightarrow \mathcal{H}_{\text {loc }}=\mathbb{C} \Omega$ (singular case)


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How to construct elements in $\mathscr{X}$
Let $A \in \mathcal{M}, B^{\prime} \in \mathcal{M}^{\prime}$, and let $L$ be a weak limit point of $\sigma_{t}\left(\alpha_{1}(A) \alpha_{-1}\left(B^{\prime}\right)\right)$ as $t \rightarrow-\infty$. Then $L \in \mathscr{X}$.

## Warping

Plan: Find a BT such that $\sigma_{t}\left(\alpha_{1}(A) \alpha_{-1}\left(B^{\prime}\right)\right) \rightarrow P_{\Omega}$ weakly as $t \rightarrow-\infty$.

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- This construction relies on a representation $T(x, y)$ of two-dimensional translation symmetry (view net $\mathcal{A}$ as chiral half of a 2d theory)
- Fix a deformation parameter $Q$, an antisymmetric $(2 \times 2)$-matrix.
- Deform smooth operators $A \in \mathcal{B}(\mathcal{H})$ on smooth vects $\Psi \in \mathcal{H}$ according to

$$
A_{Q} \Psi:=(2 \pi)^{-2} \iint_{\mathbb{R}^{2} \times \mathbb{R}^{2}} e^{-i p x} T(Q p) A T(-Q p) \cdot T(x) \Psi d p d x
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reminiscent of Weyl-Moyal product.

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## Facts [Buchholz,GL,Summers 2011]:

- $A_{Q}$ extends to a bounded operator. $A \mapsto A_{Q}$ is a faithful representation of the Rieffel-deformed $C^{*}$-algebra $\left(\mathcal{C}_{Q}, \times_{Q},\|\cdot\|_{Q}\right)$.
- Let $\mathcal{M}_{Q}:=\left\{A_{Q}: A \in \mathcal{M} \text { smooth }\right\}^{\prime \prime}$. If $\kappa \geq 0$, then also $\left(\mathcal{M}_{Q}, T, \Omega\right)$ is a Borchers triple.

Theorem ([GL/Scotford 2021] Deforming the free field leads to singular BTs)
Consider the free field triple $(\mathcal{M}, T, \Omega)$. Then there exist operators $A$ affiliated to $\mathcal{M}$ and $B^{\prime}$ affiliated to $\mathcal{M}^{\prime}$ such that for $\kappa>0$

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\underset{t \rightarrow-\infty}{\mathrm{w}-\lim \Delta^{i t} \alpha_{1}\left(A_{Q}\right) \alpha_{-1}\left(B_{-Q}^{\prime}\right) \Delta^{-i t}=P_{\Omega} .}
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Hence $\left(\mathcal{H}_{\text {loc }}\right)_{Q}=\mathbb{C} \Omega$; this gives singular BTs.

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- Here $A, B$ can be chosen as free quantum field operators.
- The proof relies on a Riemann-Lebesgue type argument

$$
\int d p_{1} \cdots d p_{n} d q^{\prime} \overline{\Phi_{n}(p)} \Psi_{n}(p) f^{+}\left(q^{\prime}\right) \overline{g^{+}\left(q^{\prime}\right)} \prod_{l=1}^{n} e^{i\left(p_{l}, Q \Lambda_{t} q^{\prime}\right)} \longrightarrow 0 \quad \text { as } t \rightarrow-\infty
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## Conjecture

Take any 2d Borchers triple with $\mathcal{H}_{\text {loc }}=\mathcal{H}$, then deform it with deformation parameter $\kappa>0$. Then $\left(\mathcal{H}_{\text {loc }}\right)_{Q}=\mathbb{C} \Omega$ (singular case).

## Outlook/Open Questions:

- This result hints at $\mathcal{H}_{\text {loc }}$ being unstable under deformations.
- What are the intrinsic properties of this example that distinguish it from case (1) and (2)?
- Are further deformation results possible? ( $\rightarrow$ quantum group symmetry)

